

METR 5223: Atmospheric Radiation

Elementary Adjustment to Radiative Equilibrium

Lecture for Spring 2009

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$$C \frac{dT}{dt} = s - \sigma T^4$$

- C is the heat capacity per unit area, e. g. $\rho H c_w$ for a layer of water of depth H .
- s is the absorbed radiation, e. g. $\frac{1}{4} S_0 (1 - \alpha_p)$
- σT^4 assumes perfect black body emission

We have studied the equilibrium case:

$$0 = s - \sigma T^4$$

Question: Using $s = \sigma T^4$, if s increases by 1%, by what percent does T increase?

Let:

$$T(t) = \bar{T} + T'(t)$$

$$s(t) = \bar{s} + s'(t)$$

\bar{T} will be defined in terms of \bar{s} .

We will assume $T' \ll \bar{T}$

$$T^4 = (\bar{T} + T')^4 = \bar{T}^4 + 4\bar{T}^3 T' + 6\bar{T}^2 T'^2 + 4\bar{T} T'^3 + T'^4$$

$$T^4 \approx \bar{T}^4 + 4\bar{T}^3 T'$$

$$s - \sigma T^4 = \bar{s} + s' - \sigma \bar{T}^4 - \sigma 4\bar{T}^3 T'$$

$$\bar{s} - \sigma \bar{T}^4 = 0$$

$$C \frac{dT'}{dt} = \bar{s} + s' - \sigma \bar{T}^4 - 4\sigma \bar{T}^3 T'$$

$T'(t)$ is the fluctuation or departure of T from the equilibrium value \bar{T} .

$$\tau \frac{dT'}{dt} + T' = T'_0$$

$$\tau \equiv \frac{C}{4\sigma \bar{T}^3}$$

$$T'_0(t) \equiv \frac{s'}{4\sigma \bar{T}^3}$$

For $\tau = 0$, $T'(t) = T'_0(t)$.

T'_0 is the *no-lag response* to the changing $s'(t)$.

$$\tau \frac{dT'}{dt} + T' = T'_0$$

Suppose $T'_0 = 0$:

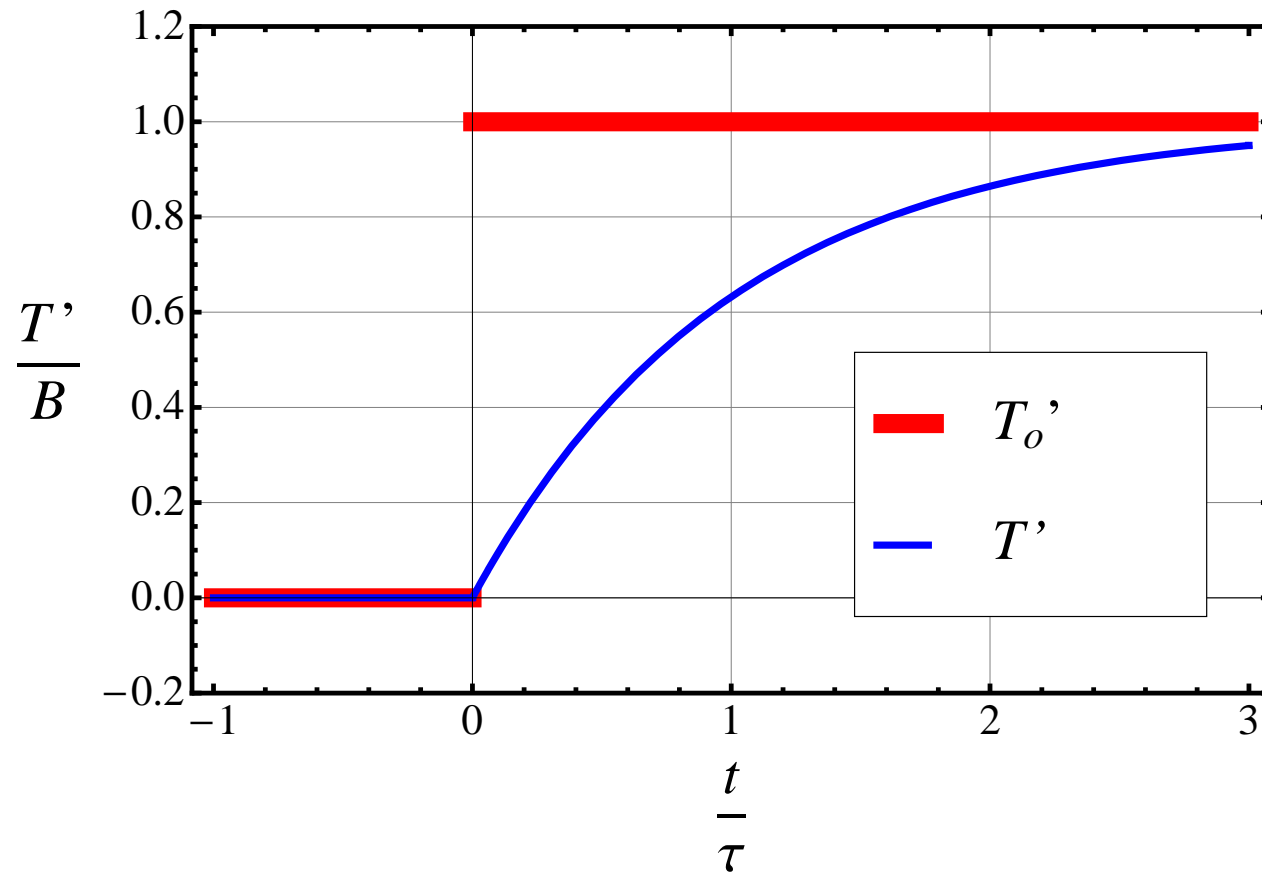
$$T'(t) = Ae^{-t/\tau}$$

Suppose $s' = 0$ for $t < 0$ and $s' = \text{constant}$ for $t \geq 0$: So we consider $T'_0 = 0$ for $t < 0$ and $T'_0 = B$ for $t > 0$. We assume $T'(0) = 0$. For $t > 0$:

$$T'(t) = Ae^{-t/\tau} + B$$

With $T'(0) = 0$:

$$T'(t) = B(1 - e^{-t/\tau})$$



$$T'(t) = B (1 - e^{-t/\tau})$$

If Earth's oceans were distributed over the entire globe, the depth would be $H = 2738$ m. With

$$\tau \equiv \frac{C}{4\sigma\bar{T}^3}$$

and

$$C = \rho H c_w,$$

$$\tau = 68 \text{ years} \quad (\bar{T} = 288K)$$

$$\tau = 97 \text{ years} \quad (\bar{T} = 255K)$$

A very important number in global warming predictions.

Instead of T'_0 changing instantly at $t = 0$, consider

$$T'_0 = Be^{t/\tau_r}$$

τ_r is the time scale for the “ramp-up” of T'_0 .

Explore how the solution depends on the ratio of the two time scales, τ/τ_r .

$$\tau \frac{dT'}{dt} + T' = Be^{t/\tau_r}$$

Propose

$$T'(t) = Ae^{t/\tau_r}$$

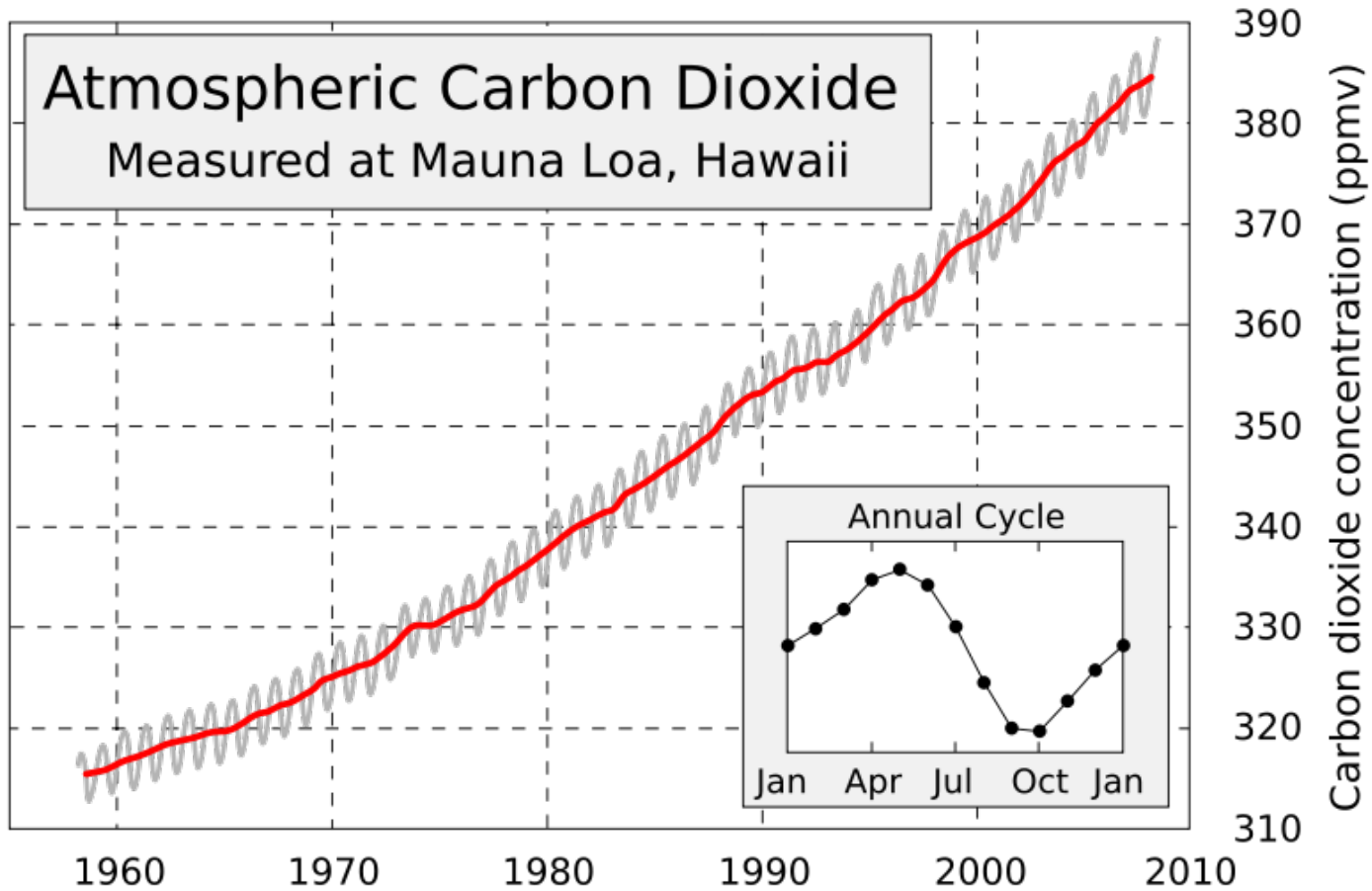
$$\tau \frac{d}{dt} A e^{t/\tau_r} + A e^{t/\tau_r} = B e^{t/\tau_r}$$

$$\frac{\tau}{\tau_r} A e^{t/\tau_r} + A e^{t/\tau_r} = B e^{t/\tau_r}$$

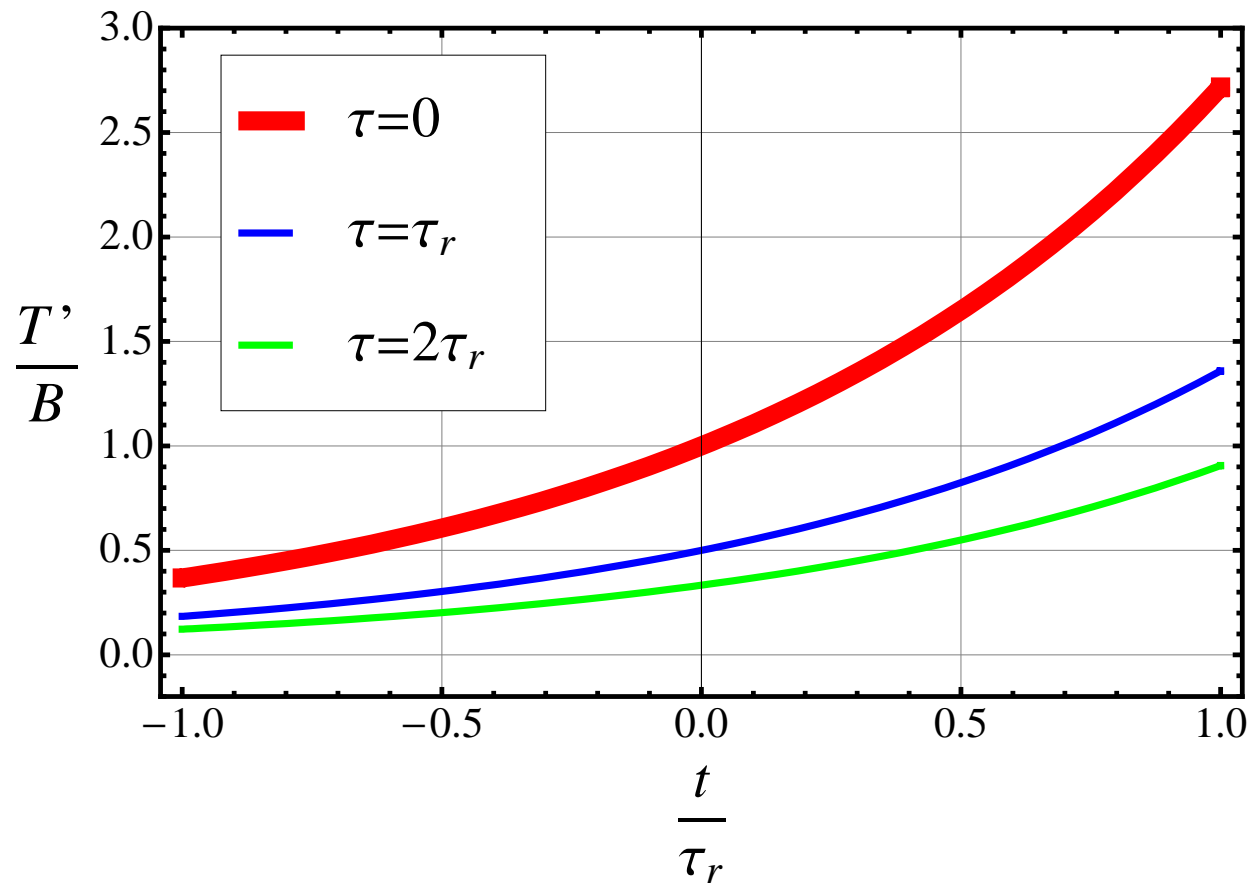
$$A \left(1 + \frac{\tau}{\tau_r} \right) = B$$

$$\frac{A}{B} = \frac{1}{1 + \frac{\tau}{\tau_r}}$$

The amplitude is diminished for larger τ/τ_r .



The increase of CO_2 from the preindustrial value of 270 ppmv is not exactly exponential. But it appears to be approximately exponential with e-folding time of about 50 years. Using 50 years for τ_r gives $A/B \approx \frac{1}{2}$.



$$T'(t) = \frac{1}{1 + \frac{\tau}{\tau_r}} B e^{-t/\tau_r}$$

Next consider seasonal forcing,

$$T'_0 = B \cos(\omega t)$$

$$\tau \frac{dT'}{dt} + T' = B \cos(\omega t)$$

To aid in comparison in the effect of the time scales, we let τ_o be the oscillation time scale, the time to go through one radian (not one cycle) so:

$$\omega = \frac{1}{\tau_o}$$

Propose

$$T'(t) = A \cos(\omega t - \delta)$$
$$\frac{dT'}{dt} = -A\omega \sin(\omega t - \delta)$$

$$-A\omega\tau \sin(\omega t - \delta) + A \cos(\omega t - \delta) = B \cos(\omega t)$$

$$- A\omega\tau [\sin(\omega t) \cos(\delta) - \cos(\omega t) \sin(\delta)]$$

$$+ A [\cos(\omega t) \cos(\delta) + \sin(\omega t) \sin(\delta)] = B \cos(\omega t)$$

What is the next step?

$$A \cos(\omega t) [\cos(\delta) + \omega\tau \sin(\delta)]$$

$$A \sin(\omega t) [-\omega\tau \cos(\delta) + \sin(\delta)] = B \cos(\omega t)$$

The only possibility for equality for all values of t and all values of $\cos(\omega t)$ and $\sin(\omega t)$ is:

$$[-\omega\tau \cos(\delta) + \sin(\delta)] = 0$$

or

$$\tan(\delta) = \omega\tau \quad \delta = \tan^{-1}(\omega\tau)$$

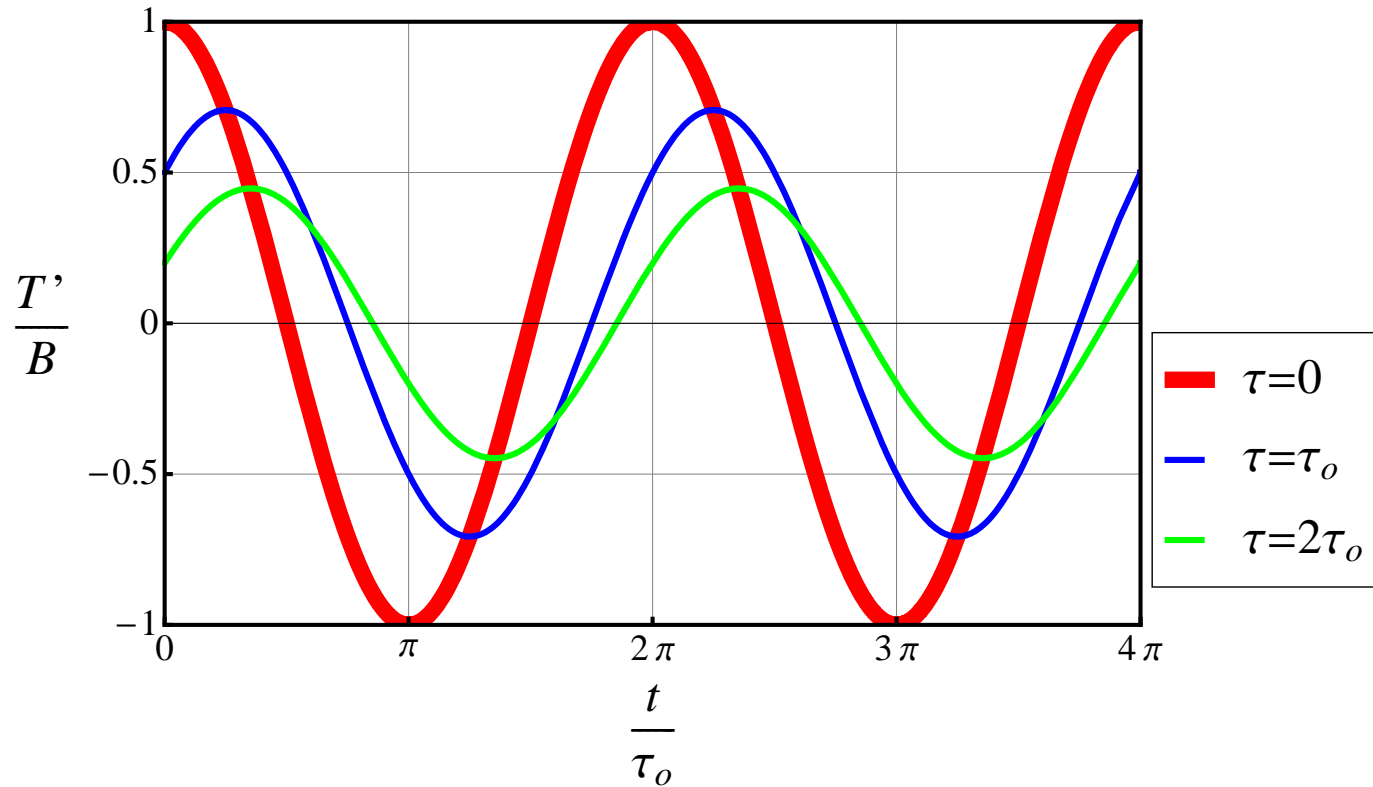
For small $\omega\tau$, $\delta \rightarrow 0$ and the response T' is in phase with the forcing T'_0 . For large $\omega\tau$, $\delta \rightarrow 90^\circ$, and the response lags the forcing by a quarter cycle.

The remaining part of the equation is:

$$A \cos(\omega t) [\cos(\delta) + \omega\tau \sin(\delta)] = B \cos(\omega t)$$

$$A [\cos(\delta) + \tan(\delta) \sin(\delta)] = B$$

$$\frac{A}{B} = \cos(\delta) = \frac{1}{\sqrt{1 + \omega^2\tau^2}} = \frac{1}{\sqrt{1 + \frac{\tau^2}{\tau_o^2}}}$$



$$T'(t) = \frac{1}{\sqrt{1 + \frac{\tau^2}{\tau_o^2}}} B \cos\left(\frac{t}{\tau_o} - \delta\right) \quad \delta = \tan^{-1}\left(\frac{\tau}{\tau_o}\right)$$

For a seasonal cycle

$$\tau_o = \frac{1 \text{ yr}}{2\pi}$$

An oceanic mixed layer of depth $H = 20$ m, gives

$\tau = 0.49$ yr:

$$\frac{\tau}{\tau_o} = 3.1$$

$$\delta = 72^\circ$$

$$A/B = \cos(\delta) = 0.3$$

In a previous lecture we studied a simple model for radiative equilibrium with a simple greenhouse. The equilibrium condition for the atmospheric slab was:

$$0 = \epsilon\sigma T_s^4 - 2\epsilon\sigma T_a^4$$

This more general energy equation for the slab would be

$$C_a \frac{dT_a}{dt} = \epsilon\sigma T_s^4 - 2\epsilon\sigma T_a^4$$

Suppose T_a is not in equilibrium, but T_s is very near equilibrium. Let

$$T_a = \bar{T}_a + T'_a(t) \quad T_s = \bar{T}_s$$
$$C_a \frac{dT'_a}{dt} = \epsilon\sigma \bar{T}_s^4 - 2\epsilon\sigma \left(\bar{T}_a^4 + 4\bar{T}_a^3 T'_a \right)$$

The equilibrium solution has the property:

$$\sigma \bar{T}_s^4 = 2\sigma \bar{T}_a^4$$

The adjustment equation simplifies to:

$$C_a \frac{dT'_a}{dt} = -8\epsilon\sigma \bar{T}_a^3 T'_a$$

The adjustment time τ is calculated with:

$$C_a = \frac{p_0}{g} c_p$$

$$\tau = \frac{C_a}{8\epsilon\sigma \bar{T}_a^3} = \frac{p_0 c_p}{g 8\epsilon\sigma \bar{T}_a^3} = 19 \text{ days}$$